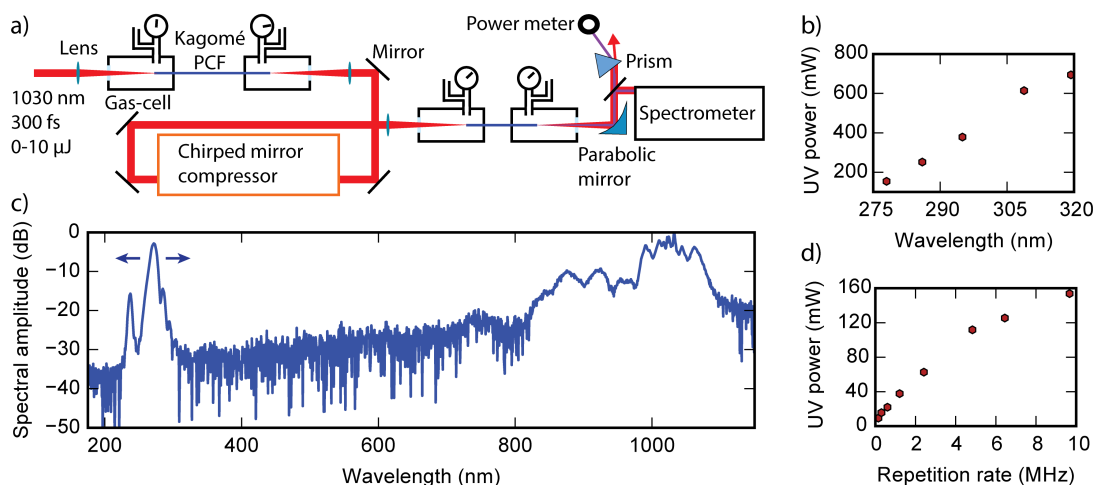


## High Average-Power and Energy Deep-Ultraviolet Femtosecond Pulse Source Driven by 10 MHz Fibre-Laser

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Fibre lasers have created a revolution in ultrafast pulse sources, due to their high average-power and high repetition-rates, with commercial 300 fs pulse systems routinely delivering  $> 20 \mu\text{J}$  and  $> 20 \text{ W}$ . However, apart from discreet harmonic schemes, short-wavelength frequency tunability is lacking, and producing shorter pulse durations requires additional compression stages. Gas-filled, hollow-core kagomé photonic-crystal fibres (kagomé-PCF) are ideal nonlinear compressors for these systems [1], and have enabled the generation of  $< 10 \text{ fs}$  pulses with  $18 \text{ W}$  average power [2]. Here we show that they can additionally be used to generate up to  $70 \text{ nJ}$  deep-UV pulses, tunable between at least  $270 \text{ nm}$  to  $320 \text{ nm}$ , at high repetition-rate, and hence average-power—exceeding  $693 \text{ mW}$ . This unique source will have a wide range of applications, to for example, angle-resolved photoemission spectroscopy [3] or ultrafast pump-probe measurements of biological molecules.



**Fig. 1** (a) The experimental set-up; (b) average deep-UV power for each pressure-tuned wavelength at  $9.6 \text{ MHz}$ ; (c) example output spectrum of the second stage when pressure tuned to produce  $278 \text{ nm}$ ; (d) corresponding average power scaling at  $278 \text{ nm}$  with repetition rate.

The experimental setup, depicted in Fig. 1(a), consists of two kagomé-PCF stages. The output of our  $1030 \text{ nm}$ ,  $300 \text{ fs}$  fibre laser is compressed to  $\sim 25 \text{ fs}$  through self-phase-modulation-based spectral broadening in a kagomé-PCF filled with a  $0$  to  $32 \text{ bar}$  krypton pressure gradient, and subsequent phase-compensation with chirped mirrors. The compressed pulses are launched into a second kagomé-PCF filled with  $10 \text{ bar}$  Ar to generate UV radiation through dispersive-wave (DW) emission [4]. By tuning the gas pressure in the second stage, and hence phase-matching conditions, the DW could be tuned, as shown in Fig. 1(b). The deep-UV power at  $9.6 \text{ MHz}$  ranges from  $153 \text{ mW}$  at  $278 \text{ nm}$ , up to  $693 \text{ mW}$  at  $320 \text{ nm}$ . For the  $278 \text{ nm}$  case the full spectrum is shown in Fig. 1(c). The deep-UV peak is broadband,  $\sim 10 \text{ nm}$ , and numerical simulations, rigorously validated in previous work [4], predict the pulse duration of the deep-UV light to be  $< 20 \text{ fs}$ . Fig. 1(d) shows how the average-power at  $278 \text{ nm}$  scales with repetition rate. The change in slope around  $6 \text{ MHz}$  is due to the onset of plasma build-up inside the fibre. This is supported by observations of side-scattered recombination luminescence (not shown), and numerical simulations.

Using lighter gases in the second stage should avoid the plasma build-up, enabling further power scaling and the generation of shorter wavelengths. Vacuum-ultraviolet DWs, down to  $120 \text{ nm}$ , were recently generated in a helium-filled kagomé-PCF system using a  $1 \text{ kHz}$ ,  $800 \text{ nm}$ , Ti:sapphire laser [5]. The results shown here raise the prospect of creating a fibre-laser pumped Watt-level VUV pulse source.

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